

Bound state solution of the Schrödinger equation for Mie potential

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Abstract

Exact solution of Schrödinger equation for the Mie potential is obtained for an arbitrary angular momentum. The energy eigenvalues and the corresponding wavefunctions are calculated by the use of the Nikiforov-Uvarov method. Wavefunctions are expressed in terms of Jacobi polynomials. The bound states are calculated numerically for some values of ℓ and n with $n \leq 5$. They are applied to several diatomic molecules.

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1 Introduction

The study of three-dimensional anharmonic oscillators has raised a considerable amount of interest due to its wide applications in molecular physics [1-4]. The Morse potential is commonly used for anharmonic oscillator. However, its wavefunction does not vanish at the origin. But those of Mie type and pseudo-harmonic potentials do. On the other hand, the Mie type potential possess the general features of the true interaction energy [1]. The Mie type and pseudo-harmonic are two special kinds of exactly solvable power-law and inverse power-law potentials other than the Coulombic and harmonic oscillator. The Hamiltonians with inverse power potential $1/r^n$ are studied in different fields. For instance, path integral, $1/N$ expansion and super-symmetrical approaches are also studied for this type of potentials [5-8]. Moreover, there are some applications of Mie type potentials such as equation of bulk metallic glass for $Zr_{48}Nb_8Cu_{14}Ni_{12}Be_{18}$ in solid state physics [9], and other applications in solid state physics [10], molecular spectroscopy [11], fluid mechanics [12], the interatomic interaction potential in molecular physics [13], $1/r^n$ potential in three dimensions [14].

The Nikiforov-Uvarov (NU) method [15] is introduced for the solution of the hypergeometric type second order differential equations [16] appeared in the time-independent problems. Recently, NU-method has been used to solve Schrödinger equation for some well known potentials [17-22], Dirac and Klein-Gordon equations for the Coulomb and some exponential potentials. In the present work the Schrödinger equation is solved by the NU method for Mie potential with any value of angular momentum ℓ . As an illustration, energy eigenvalues are computed for N_2 , CO , NO and CH molecules.

The organization of the paper is as follows. In section II, NU method is briefly introduced. Mie potential calculations are done by using NU method in section III. Results are discussed in section IV.

2 The Nikiforov-Uvarov Method

The NU-method developed by Nikiforov and Uvarov (NU-method) is based on reducing the second order differential equations (ODEs) to a generalized equation of hypergeometric type [15]. It is provided us an analytic solution of Schrödinger equation for certain kind of potentials. It is based

on the solutions of general second order linear differential equation with orthogonal function [16]. The one-dimensional Schrödinger equation is reduced to a generalized equation of hypergeometric type with an appropriate coordinate transformation $s = s(r)$. Then, the equation has the form,

$$\psi''(s) + \frac{\tilde{\tau}(s)}{\sigma(s)}\psi'(s) + \frac{\tilde{\sigma}(s)}{\sigma^2(s)}\psi(s) = 0 \quad (1)$$

where $A(s) = \frac{\tilde{\tau}(s)}{\sigma(s)}$ and $B(s) = \frac{\tilde{\sigma}(s)}{\sigma^2(s)}$. In Eq. (1), $\sigma(s)$ and $\tilde{\sigma}(s)$ are polynomials at most second degree, and $\tilde{\tau}(s)$ is a polynomial with at most first degree [15]. The wave function is constructed as a multiple of two independent parts,

$$\psi(s) = \phi(s)y(s), \quad (2)$$

and Eq. (1) becomes

$$\sigma(s)y''(s) + \tau(s)y'(s) + \lambda y(s) = 0, \quad (3)$$

where

$$\sigma(s) = \pi(s)\frac{d}{ds}(\ln\phi(s)), \quad (4)$$

and

$$\tau(s) = \tilde{\tau}(s) + 2\pi(s). \quad (5)$$

λ is defined as

$$\lambda_n + n\tau' + \frac{[n(n-1)\sigma'']}{2} = 0, n = 0, 1, 2, \dots \quad (6)$$

We determine $\pi(s)$ and λ by defining

$$k = \lambda - \pi'(s). \quad (7)$$

From the Eqs. (4-5), $\pi(s)$ becomes

$$\pi(s) = \left(\frac{\sigma' - \tilde{\tau}}{2}\right) \pm \sqrt{\left(\frac{\sigma' - \tilde{\tau}}{2}\right)^2 - \tilde{\sigma} + k\sigma} \quad (8)$$

$\pi(s)$ has to be a polynomial of degree at most one in the Eq. (8), so the expression under the square root must be the square of a polynomial of first degree [15]. This is simply possible only if its discriminant is zero. After defining k , one can obtain $\pi(s)$, $\tau(s)$, $\phi(s)$ and λ . If we look at the Eq. (4) and the Rodrigues relation

$$y_n(s) = \frac{B_n}{\rho(s)} \frac{d^n}{ds^n} [\sigma^n(s) \rho(s)], \quad (9)$$

where C_n is normalizable constant and the weight function satisfy the relation as

$$\frac{d}{ds} [\sigma(s) \rho(s)] = \tau(s) \rho(s). \quad (10)$$

where

$$\frac{\phi'(s)}{\phi(s)} = \frac{\pi(s)}{\sigma(s)}. \quad (11)$$

3 Mie Potential Calculations

The Mie-type potentials are given by

$$V(r) = \epsilon \left[\frac{k}{\ell - k} \left(\frac{a}{r}\right)^\ell - \frac{\ell}{\ell - k} \left(\frac{a}{r}\right)^k \right] \quad (12)$$

where ϵ is the interaction energy between two atoms in a solid at $x = a$, and $\ell > k$ is always satisfied.

We solve the one-dimensional Mie potential [1-2] with $\ell = 2k$ combination. By choosing the special case $k = 1$, corresponding to a Coulombic-type potential with an additional centrifugal potential barrier, we get the following form:

$$V(r) = V_0 \left[\frac{1}{2} \left(\frac{a}{r} \right)^2 - \left(\frac{a}{r} \right) \right]; \quad V_0 = 2\epsilon k \quad (13)$$

where V_0 is the dissociation energy and a is the positive constant which is strongly repulsive at shorter distances. The radial part of the Schrödinger equation for a diatomic molecule potential is

$$\left[-\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d}{dr} \right) + \frac{\ell(\ell+1)\hbar^2}{2\mu r^2} + V(r) \right] R_{n\ell}(r) = E_{n\ell} R_{n\ell}(r), \quad (14)$$

where μ is the reduced mass of the diatomic molecules. n denotes the radial quantum number (n and ℓ are named as the vibration-rotation quantum numbers in molecular chemistry). r is the internuclear separation. Substituting the explicit form of $V(r)$, we get

$$\frac{d^2 R_{n\ell}(r)}{dr^2} + \frac{2}{r} \frac{dR_{n\ell}(r)}{dr} + \frac{2\mu}{\hbar^2} \left[E_{n\ell} - V_0 \left(\frac{1}{2} \left(\frac{a}{r} \right)^2 - \frac{a}{r} \right) \right] R_{n\ell}(r) - \frac{\ell(\ell+1)\hbar^2}{2\mu r^2} R_{n\ell}(r) = 0. \quad (15)$$

By defining the following variables:

$$\varepsilon^2 = E \quad (16)$$

$$\beta = -\frac{2\mu}{\hbar^2} V_0 a \quad (17)$$

$$\gamma = \frac{2\mu}{\hbar^2} \left(\frac{1}{2} V_0 a^2 + \frac{\ell(\ell+1)\hbar^2}{2\mu} \right) \quad (18)$$

The Schrödinger equation takes the simple form:

$$\left[\frac{d^2}{dr^2} + \frac{2}{r} \frac{d}{dr} + \frac{1}{r^2} (\varepsilon^2 r^2 - \beta r - \gamma) \right] R_{n\ell}(r) = 0 \quad (19)$$

From Eq. (19), it is clear that $\sigma(r) = r$, $\tilde{\tau} = 2$, $\tilde{\sigma} = \varepsilon^2 r^2 - \beta r - \gamma$. We find $\pi(r)$ as

$$\pi(r) = -\frac{1}{2} \pm \begin{cases} i\varepsilon r + \frac{1}{2}\sqrt{4\gamma+1}, & k_1 = -\beta + \varepsilon\sqrt{-1-4\gamma}; \\ i\varepsilon r - \frac{1}{2}\sqrt{4\gamma+1}, & k_2 = -\beta - \varepsilon\sqrt{-1-4\gamma}. \end{cases} \quad (20)$$

and $\tau(r)$ can be written as

$$\tau(r) = \begin{cases} 1 + 2i\varepsilon r + \sqrt{4\gamma+1}, & k_1 = -\beta + \varepsilon\sqrt{-1-4\gamma}; \\ 1 - 2i\varepsilon r + \sqrt{4\gamma+1}, & k_2 = -\beta - \varepsilon\sqrt{-1-4\gamma}. \end{cases} \quad (21)$$

For appropriate solutions $\tau'(r) < 0$ [15], we use $k_2 = -\beta - \varepsilon\sqrt{-1 - 4\gamma}$, $\pi_2 = -\frac{1}{2} - i\varepsilon r - \frac{1}{2}\sqrt{1 + 4\gamma}$ and $\tau(r) = 1 - 2i\varepsilon r + \sqrt{1 + 4\gamma}$ for obtaining eigenvalues and eigenfunctions. Using Eqs. (6-7) we find energy spectrum as

$$E_n = -\left(\frac{2\mu V_0 a}{\hbar^2}\right)^2 \left[2n + 1 + \sqrt{1 + \frac{8\mu}{\hbar^2} \left(\frac{V_0 a^2}{2} + \frac{\ell(\ell + 1)\hbar^2}{2\mu}\right)}\right]^{-2} \quad (22)$$

Wavefunction is calculated from Eqs. (9-11).

$$\phi(r) = r^{\frac{1}{2}(-1+\sqrt{1+4\gamma})} e^{-i\varepsilon r} \quad (23)$$

and

$$\rho(r) = r^{\sqrt{1+4\gamma}} e^{-2i\varepsilon r} \quad (24)$$

and the radial part of the wave function becomes

$$R_{n,\ell} = C_n r^{\frac{1}{2}(-1+\sqrt{1+4\gamma})} e^{-i\varepsilon r} L_n^{\sqrt{1+4\gamma}}(2i\varepsilon r) \quad (25)$$

If we rewrite Eq. (22) in atomic units, we obtain

$$E_n = -V_0^2 a^2 \left(2n + 1 + \sqrt{(2\ell + 1)^2 + 2V_0 a^2}\right)^{-2} \quad (26)$$

4 Conclusions

We have investigated the analytical solution of the Schrödinger equation for Mie potential. By using a special case of Mie potential as $m = 2k$ and $k = 1$, the problem is reduced to a Coulombic potential with the additional centrifugal barrier. Energy eigenvalues and corresponding eigenfunctions are obtained by the NU method. energy eigenvalues are computed for N_2 , CO , NO and CH molecules. They are listed in Table 1.

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Table 1: Calculated energy eigenvalues of the Mie potential for N_2 , CO , NO and CH diatomic molecules with different values of n and ℓ in eV . The data for potential parameters are taken from [22].

State (n)	ℓ	N_2	CO	NO	CH
0	0	0.00271631	0.0254063	0.0205538	0.0416018
1	0	0.0810235	0.0756383	0.0611503	0.120556
	1	0.0810268	0.0758723	0.0613568	0.122185
2	0	0.13411	0.125172	0.101132	0.194768
	1	0.134351	0.125403	0.101335	0.1963
	2	0.134833	0.125867	0.101742	0.199356
3	0	0.186486	0.17402	0.140511	0.264609
	1	0.186719	0.174248	0.140712	0.266052
	2	0.187194	0.174704	0.141112	0.26893
	3	0.187907	0.175387	0.141713	0.273228
4	0	0.238152	0.222196	0.1793	0.330417
	1	0.238385	0.222421	0.179498	0.331777
	2	0.238854	0.22287	0.179892	0.334491
	3	0.239557	0.223544	0.180484	0.338544
	4	0.240495	0.224442	0.181272	0.343916
5	0	0.28913	0.269712	0.217511	0.392495
	1	0.289361	0.269933	0.217705	0.393779
	2	0.289824	0.270377	0.218094	0.396341
	3	0.290518	0.271041	0.218677	0.400166
	4	0.291443	0.271927	0.219453	0.405238
	5	0.292599	0.273036	0.220424	0.41153